Technical Notes

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Coupled Radiation with Turbulent Convection in Electric Arcs

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Introduction

THE satisfactory performance of plasma arc devices has been known to be limited by Joule heating, as well as thermal conduction, turbulent convection and radiation. The effect of turbulence has not been included in previous investigations, and this may restrict the validity of the reported results to arcs operating with small flow rates. The present paper accounts for the coupling of turbulence with radiation in arcs, with results shown to indicate the relative role of each at various values of Reynolds number and operating current.

Analysis

Consider a steady flow in a cooled, constant wall-temperature tube whose internal surface is black. The body forces and viscous dissipation are negligible. The plasma is assumed to be in a state of local thermodynamic equilibrium. The axial field intensity is constant, and diffusion of ions and electrons is ambipolar.

The governing equations in the fully developed stage are

$$F = -(\mu + \rho \varepsilon_{M}) \, du/dr \tag{1}$$

$$(1/r) d/dr(rq) = \delta E^2$$
 (2)

$$q = q_R - (\lambda + \rho c_p \varepsilon_H) dT/dr \tag{3}$$

$$q_c = -\lambda \, dT/dr \tag{4}$$

Where c_p is the specific heat at constant pressure, E is the field intensity, F is the shear stress, q is the total heat flux, q_c is the conduction heat flux, q_R is the total radiation heat flux, r is the radial coordinate, T is the absolute temperature, u is the velocity, ε_M and ε_H are the turbulent diffusivities for momentum and heat, respectively, λ is the thermal conductivity, δ is the electrical conductivity, μ is the molecular viscosity, and ρ is the density. The boundary conditions are

i)
$$u = 0$$
, $r = r_0$ ii) $q = 0$, $r = 0$ iii) $T = T_0$, $r = r_0$

the suffix 0 used here and later on refers to the quantities evaluated at the wall. The assumption that turbulence has direct effect only on viscosity and thermal conductivity has been justified by Harder and Cann.¹

The general expression for the radiation heat flux taking into

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account the spectral dependence of the absorption coefficient is given by

$$q_{R}(r) = \frac{4}{\pi} \int_{0}^{\pi/2} \cos(\gamma) \left\{ \int_{r\sin(\gamma)}^{r_{0}} \frac{de}{dr'} \left(-\varepsilon \left[T(r'), \frac{1}{r} - \frac{1}{r} \sin(\gamma) \right] \right) dr' \right\} d\gamma$$

$$(r + r' - 2r\sin(\gamma)) / \cos(\gamma) + \varepsilon \left[T(r'), \frac{1}{r} - \frac{1}{r} \cos(\gamma) \right] dr' \right\} d\gamma$$
(5)

where

$$\varepsilon[T(r'), X] = \int_0^\infty \frac{de_\omega(r')}{de(r')} (1 - D_3[K_\omega(r')X]) d\omega$$

$$D_n(X) = \int_0^{\pi/2} \cos(\alpha) \exp(-X/\cos(\alpha)) d\alpha \qquad (6)$$

Here e is the black body emissive power ($=\sigma T^4$, σ is called Stefan-Boltzmann's constant), e_{ω} is the spectral emissive power, K_{ω} is the monochromatic absorption coefficient, and ε is called the modified emissivity. The details of the derivation of Eq. (5) are given by Pyare and Abu-Romia.²

In accordance with Ref. 3, eddy diffusivity is taken as

$$\varepsilon_{M} = k_{1}^{2} uy \left[1 - \exp\left(-\frac{k_{1}^{2} uy}{\mu/\rho}\right) \right] \text{ for } \dot{y}^{+} \leq 26$$
 (7)

where $k_1 = 0.124$. For $y^+ > 26$

$$\varepsilon_M = k_2 (F_0/\rho)^{1/2} y (1 - y/r_0) - \mu/\rho$$
 (8)

assuming linear shear stress distribution, and

$$\varepsilon_{\rm M} = k_2 (K_0/\rho)^{1/2} y - \mu/\rho \tag{9}$$

assuming constant shear stress distribution. Here $k_2=0.36$, $y=r_0-r$ and $y^+=(F_0/\rho_0)^{1/2}/(\mu_0/\rho_0)$. The equality of eddy diffusivities is postulated in the analysis, i.e., $\varepsilon_M=\varepsilon_H$.

In the computations, the transport and thermodynamic properties were taken from Yos⁴ and the monochromatic absorption coefficient from Churchill et al.⁵ The results were obtained by assuming a polynomial for the fourth power of the temperature,

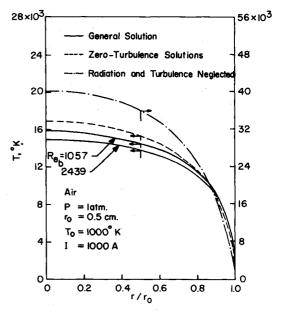


Fig. 1 Effect of turbulence on temperature profiles.

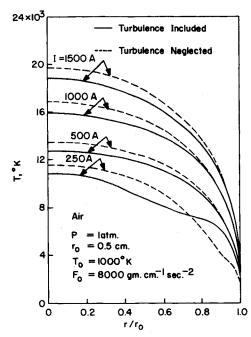


Fig. 2 Temperature profiles for various values of current.

with the coefficients being obtained by an iterative scheme. Details of the numerical results are given in Ref. 6. The current *I*, indicated in Figs. 1–3, is related to the voltage by the over-all energy balance

$$2\pi r_0 q_0 = EI \tag{10}$$

Values of Reynolds number Re_b , mass flow rate m, and the mean bulk enthalpy are calculated from

$$Re_b = 2m/(\pi r_0 \mu_b) \tag{11}$$

$$m = 2\pi \int_0^{r_0} \rho u r \, dr \tag{12}$$

$$H_b = \left(2\pi \int_0^{r_0} H\rho u r \, dr\right) / m \tag{14}$$

where μ_h is the viscosity evaluated at the bulk temperature.

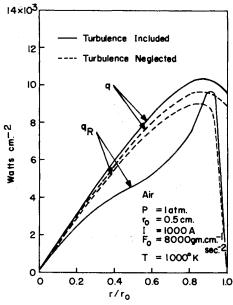


Fig. 3 Effect of turbulence on heat flux profiles.

Results

As indicated in Fig. 1, the effect of turbulence, even for a relatively high Reynolds number, is small compared to the effect of radiation which is the dominant mechanism at high currents. At lower operating currents, i.e., low plasma temperatures, the effect of turbulence is more significant while radiation has a diminishing role. As can be noted in Fig. 2, the temperature profile with turbulence is flatter than that without turbulence. Fig. 3 shows that turbulence has significant effect on radiation flux distribution, and as expected, the total heat flux at the wall is increased by the presence of turbulence.

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Dynamically-Induced Large Deformations of Multilayer, Variable Thickness Shells

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In recent years the analysis and prediction of large transient and permanent deformations of structures have been of increasing interest in connection with explosive and/or magnetic forming of simple and complex structural shapes. Also various aerospace, water, and land vehicles as well as stationary structures may be subject to intense transient loading caused by

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